

## THE FLOW-CURVATURE OF PLANE PARAMETRIZED CURVES

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ABSTRACT. We introduce and study a new frame and a new curvature function for a fixed parametrization of a plane curve. This new frame is called *flow* since it involves the time-dependent rotation of the usual Frenet flow; the angle of rotation is exactly the current parameter. The flow-curvature is calculated for several examples obtaining the logarithmic spirals (and the circle as limit case) and the Grim Reaper as flat-flow curves. A main result is that the scaling with  $\frac{1}{\sqrt{2}}$  of both Frenet and flow-frame belong to the same fiber of the Hopf bundle. Moreover, the flow-Fermi-Walker derivative is defined and studied.

The theory of geometric flows is a new and fascinating field of research in geometric analysis. The most simple of them is *the curve shortening flow* and already the excellent survey [3] is twenty years old. Recall that the main geometric tool in this last flow is the well-known curvature of plane curves. Hence, to give a re-start to this problem seems to search for variants of the curvature, or in terms of [8], *deformations* of the usual curvature. The goal of this short note is to propose such a deformation which in turn defines a Fermi-Walker type derivative.

Fix an open interval  $I \subseteq \mathbb{R}$  and consider  $C \subset \mathbb{R}^2$  a regular parametrized curve of equation:

$$C : r(t) = (x(t), y(t)) = x(t)\bar{i} + y(t)\bar{j}, \quad \|r'(t)\| > 0, \quad t \in I. \quad (1)$$

The ambient setting  $\mathbb{R}^2$  is an Euclidean vector space with respect to the canonical inner product:

$$\langle u, v \rangle = x^1 y^1 + x^2 y^2, \quad u = (x^1, x^2) \in \mathbb{R}^2, \quad v = (y^1, y^2) \in \mathbb{R}^2, \quad 0 \leq \|u\|^2 = \langle u, u \rangle. \quad (2)$$

The infinitesimal generator of the rotations in  $\mathbb{R}^2 = \mathbb{C}$  is the linear vector field, called *angular*:

$$\xi(u) := -x^2 \frac{\partial}{\partial x^1} + x^1 \frac{\partial}{\partial x^2}, \quad \xi(u) = i \cdot u = i \cdot (x^1 + ix^2), \quad i = \sqrt{-1}. \quad (3)$$

It is a complete vector field with integral curves the circles  $\mathcal{C}(O, r)$ :

$$\begin{cases} \gamma_{u_0}^\xi(t) = (u_0^1 \cos t - u_0^2 \sin t, u_0^1 \sin t + u_0^2 \cos t) = R(t) \cdot \begin{pmatrix} u_0^1 \\ u_0^2 \end{pmatrix}, & t \in \mathbb{R}, \\ r = \|u_0\| = \|(u_0^1, u_0^2)\|, & R(t) := \begin{pmatrix} \cos t & -\sin t \\ \sin t & \cos t \end{pmatrix} \in SO(2) = S^1 \end{cases} \quad (4)$$

and since the rotations  $R(t)$  are isometries of the Riemannian metric  $g_{can} = dx^2 + dy^2 = |dz|$  it follows that  $\xi$  is a Killing vector field of the Riemannian manifold  $(\mathbb{R}^2, g_{can})$ . The first integrals of  $\xi$  are the Gaussian functions i.e. multiples of the square norm:  $f_\alpha(x, y) = \alpha(x^2 + y^2)$ ,  $\alpha \in \mathbb{R}$ . For an arbitrary vector field  $X = A(x, y)\frac{\partial}{\partial x} + B(x, y)\frac{\partial}{\partial y}$  its Lie bracket with  $\xi$  is:

$$[X, \xi] = (yA_x - xA_y - B)\frac{\partial}{\partial x} + (A + yB_x - xB_y)\frac{\partial}{\partial y}, \quad (5)$$

where the subscript denotes the variable corresponding to the partial derivative. For example,  $\xi$  commutes with *the radial* (or Euler) vector field  $E(x, y) = x\frac{\partial}{\partial x} + y\frac{\partial}{\partial y}$ , which is also a complete vector field having as integral curves the homotheties  $\gamma_{u_0}^E(t) = e^t u_0$  for all  $t \in \mathbb{R}$ .

The Frenet apparatus of the curve  $C$  is provided by:

$$\begin{cases} T(t) = \frac{r'(t)}{\|r'(t)\|}, & N(t) = i \cdot T(t) = \frac{1}{\|r'(t)\|}(-y'(t), x'(t)), \\ k(t) = \frac{1}{\|r'(t)\|} \langle T'(t), N(t) \rangle = \frac{1}{\|r'(t)\|^3} \langle r''(t), ir'(t) \rangle = \frac{1}{\|r'(t)\|^3} [x'(t)y''(t) - y'(t)x''(t)]. \end{cases} \quad (6)$$

Hence, if  $C$  is naturally parametrized (or parametrized by arc-length) i.e.  $\|r'(s)\| = 1$  for all  $s \in I$  then  $r''(s) = k(s)ir'(s)$ . In a complex approach based on  $z(t) = x(t) + iy(t) \in \mathbb{C} = \mathbb{R}^2$  we have:

$$\begin{cases} k(t) = \frac{1}{|z'(t)|^3} \text{Im}(\bar{z}'(t) \cdot z''(t)) = \frac{1}{|z'(t)|} \text{Im}\left(\frac{z''(t)}{z'(t)}\right) = \frac{1}{|z'(t)|} \text{Im}\left[\frac{d}{dt}(\ln z'(t))\right] \in \mathbb{R}, \\ \text{Re}(\bar{z}'(t) \cdot z''(t)) = \frac{1}{2} \frac{d}{dt} \|r'(t)\|^2, & f_\alpha(z) = \alpha|z|^2. \end{cases} \quad (7)$$

The multiplication with the complex unit  $i$  corresponds to the rotation  $R\left(\frac{\pi}{2}\right)$ ; we have also:

$$\frac{d}{dt} R(t) = R\left(t + \frac{\pi}{2}\right) = R(t)R\left(\frac{\pi}{2}\right) = R\left(\frac{\pi}{2}\right)R(t), \quad (8)$$

and the Frenet equations can be unified by means of the column matrix  $\mathcal{F}(t) = \begin{pmatrix} T \\ N \end{pmatrix}(t)$  as:

$$\frac{d}{dt} \mathcal{F}(t) = \|r'(t)\| k(t) R\left(-\frac{\pi}{2}\right) \mathcal{F}(t). \quad (9)$$

It is an amazing fact that if the general rotation  $R(t)$  belongs to the Lie group  $SO(2) = S^1$  its particular values  $R\left(\pm\frac{\pi}{2}\right)$  are elements of its Lie algebra  $so(2)$  of skew-symmetric  $2 \times 2$  matrices. In fact,  $\{R\left(\frac{\pi}{2}\right)\}$  is exactly the basis of  $so(2)$ .

This short note defines a new frame and correspondingly a new curvature function for  $C$ :

**Definition 1.** *The flow-frame of  $C$  consists in the pair of unit vectors  $(E_1^f(t), E_2^f(t)) \in T^2 := S^1 \times S^1$  given by:*

$$\mathcal{E}(t) := \begin{pmatrix} E_1^f \\ E_2^f \end{pmatrix} (t) = R(t)\mathcal{F}(t) = \begin{pmatrix} \cos tT(t) - \sin tN(t) \\ \sin tT(t) + \cos tN(t) \end{pmatrix} \quad (10)$$

the letter  $f$  being the initial of the word "flow". The flow-curvature of  $C$  is the smooth function  $k_f : I \rightarrow \mathbb{R}$  given by the flow-equations:

$$\frac{d}{dt}\mathcal{E}(t) = \|r'(t)\|k_f(t)R\left(-\frac{\pi}{2}\right)\mathcal{E}(t). \quad (11)$$

Before starting its study we point out that this work is dedicated the memory of Academician Radu Miron (1927-2022). He was always interested in the geometry of curves and besides his theory of *Myller configurations* ([10]) he generalized also a type of curvature for space curves in [9]. We remark also that this note follows the idea of Bishop in his delightful note [2] and that the flow-curvature of spacelike parametrized curves in the Lorentz plane was introduced by the author in [4].

Returning to our subject we note as a first main result:

**Proposition 1.** *The expression of the flow-curvature is:*

$$k_f(t) = k(t) - \frac{1}{\|r'(t)\|} < k(t). \quad (12)$$

As a consequence, the curve  $C$  and its trigonometrical rotation  $iC$  share the same flow-curvature.

**Proof** We have directly in the flow-frame:

$$\|r'(t)\|k_f(t)R\left(-\frac{\pi}{2}\right) = R\left(t + \frac{\pi}{2}\right)R(-t) + \|r'(t)\|k(t)R(t)R\left(-\frac{\pi}{2}\right)R(-t) \quad (13)$$

and the conclusion follows. Concerning the consequence it is obvious that  $C$  and  $iC : t \rightarrow (-y(t), x(t))$  share the same curvature  $k$  and the same second term from (12).  $\square$

**Example 1.** *i) If  $C$  is the line  $r_0 + tu, t \in \mathbb{R}$  with the vector  $u \neq \bar{0} = (0, 0)$  then  $k_f$  is constant:*

$$k_f(t) = -\frac{1}{\|u\|} = \text{constant} < 0. \quad (14)$$

*In particular, if  $u$  is an unit vector then  $k_f(t) = -1$ .*

*ii) The circle  $\mathcal{C}(O, R)$  with the usual parametrization  $r(t) = Re^{it}$  is a flat-flow curve i.e.  $k_f = 0$ . Indeed, the flow-frame is constant and universal for the families of*

concentric circles i.e. it does not depend on the radius  $R$  (exactly as the Frenet frame):

$$E_1^f = (0, 1) = \bar{j}, \quad E_2^f = (-1, 0) = -\bar{i}. \quad (15)$$

More generally, if  $C$  is expressed in polar coordinates as  $C : \rho = \rho(t)$  for  $t \in I$  then  $C$  is a flat-flow curve if and only if  $C$  is a logarithmic spiral  $\rho_{R,\alpha}(t) = Re^{\alpha t}$ ,  $R, \alpha > 0$  and  $t \in \mathbb{R}$ . The limit case  $\alpha \rightarrow 0$  gives the circle  $C(O, R)$  and the flow-frame of the logarithmic spiral is:  $E_1^f = \frac{1}{\sqrt{\alpha^2+1}}(\alpha, 1)$ ,  $E_2^f = \frac{1}{\sqrt{\alpha^2+1}}(-1, \alpha)$ ; if  $\alpha = \cot \varphi$  then  $E_1^f = e^{\varphi i}$ ,  $E_2^f = e^{i(\frac{\pi}{2}+\varphi)}$ .

iii) Fix  $R \in (0, +\infty)$  and the plane curve  $C : w = F(Re^{it})$  with  $t$  as an increasing parameter and  $F = F(z)$  a holomorphic function. Then the curvatures are:

$$k(t) = \frac{1}{|zF'(z)|} \operatorname{Re} \left( 1 + \frac{zF''(z)}{F'(z)} \right), \quad k_f(t) = \frac{1}{|zF'(z)|} \operatorname{Re} \left( \frac{zF''(z)}{F'(z)} \right). \quad (16)$$

For the circle example of  $F(z) = z^2$  it results  $k = \frac{1}{R^2} = \text{constant}$  and  $k_f = \frac{1}{2R^2} = \text{constant}$  which proves the proper dependence of  $k_f$  on the parametrizations of  $C$ .  $\square$

**Remark 1.** i) Suppose that  $I$  is symmetric with respect to  $0 \in \mathbb{R}$  and that  $C$  is positively oriented in the terms of Definition 1.14 from [13, p. 17]. Suppose also the  $C$  is convex; then applying the Theorem 1.18 of page 19 from the same book it results for the usual curvature the inequality  $k \geq 0$ . Hence the opposite curve  $C^- : t \in I \rightarrow r(-t)$  has the flow-curvature  $k_f < 0$ .

ii) An important tool in dynamics is the Fermi-Walker derivative. Let  $\mathcal{X}_C$  be the set of vector fields along the curve  $C$ . Then the Fermi-Walker derivative is the map ([5])  $\nabla_C^{FW} : \mathcal{X}_C \rightarrow \mathcal{X}_C$ :

$$\nabla_C^{FW}(X) := \frac{d}{dt}X + \|r'(\cdot)\|k[\langle X, N \rangle T - \langle X, T \rangle N] = \frac{d}{dt}X + \|r'(\cdot)\|k[X^\flat(N)T - X^\flat(T)N] \quad (17)$$

with  $X^\flat$  the differential 1-form dual to  $X$  with respect to the Euclidean metric. In a matrix form we can express this as follows:

$$\nabla_C^{FW} = \frac{d}{dt} - \|r'\|k \begin{vmatrix} (\cdot)^\flat(T) & (\cdot)^\flat(N) \\ T & N \end{vmatrix} = \frac{d}{dt} + \|r'\|k \begin{vmatrix} T & (\cdot)^\flat(T) \\ N & (\cdot)^\flat(N) \end{vmatrix}. \quad (18)$$

It is natural to make here a remark concerning rotation-minimizing fields  $X \in \mathcal{X}_C$  i.e. fields satisfying:

$$\frac{d}{dt}X(t) = \lambda(t)T(t), \quad \langle X(t), T(t) \rangle = 0$$

for a smooth function  $\lambda = \lambda(t)$ . Then the Fermi-Walker derivative of such  $X$  is also parallel with the tangent  $T$ :

$$\nabla_C^{FW} X(t) = [\lambda(t) + \|r'(t)\|k(t)\langle X(t), N(t) \rangle]T(t).$$

Calculating the Fermi-Walker derivative on our frames we get:

$$\nabla_C^{FW}(T) = \nabla_C^{FW}(N) = 0, \quad \nabla_C^{FW}(E_1^f) = -E_2^f, \quad \nabla_C^{FW}(E_2^f) = E_1^f. \quad (19)$$

With the matrix notation we can express these relations as:

$$\nabla_C^{FW}(\mathcal{F}) = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \quad \nabla_C^{FW}(\mathcal{E}) = R\left(\frac{\pi}{2}\right)\mathcal{E} \quad (20)$$

and the Fermi-Walker derivative can be expressed in terms of  $k_f$  as:

$$\nabla_C^{FW}(X) = \frac{d}{dt}X + (1 + \|r'\|k_f)[X^b(N)T - X^b(T)N]. \quad (21)$$

Also, we can define the flow-Fermi-Walker derivative as:

$$\nabla_C^{fFW}(X) := \frac{d}{dt}X + \|r'(\cdot)\|k_f[X^b(N)T - X^b(T)N] = \nabla_C^{FW}(X) + T \wedge N(X) \quad (22)$$

with the skew-symmetric endomorphism  $\wedge \in so(2)$  defined by:

$$X \wedge Y := \langle X, \cdot \rangle Y - \langle Y, \cdot \rangle X = (X^1Y^2 - X^2Y^1)R\left(\frac{\pi}{2}\right), \quad X = (X^1, X^2), \quad Y = (Y^1, Y^2).$$

Then:

$$\nabla_C^{fFW}(\mathcal{F}) = R\left(-\frac{\pi}{2}\right)\mathcal{F}, \quad \nabla_C^{fFW}(\mathcal{E}) = \begin{pmatrix} 0 \\ 0 \end{pmatrix}. \quad (23)$$

As in the usual case, if  $V, W \in \mathcal{X}_C$  are flow-Fermi-Walker fields i.e. with zero flow-Fermi-Walker derivative then the value  $\langle V, W \rangle \in \mathbb{R}$  is constant along  $C$ .

iii) Remark that the 4-dimensional vectors  $\frac{1}{\sqrt{2}}\mathcal{F}$  and  $\frac{1}{\sqrt{2}}\mathcal{E}$  belong to the Clifford torus  $\frac{1}{\sqrt{2}}T^2 \subset S^3$ . A remarkable Riemannian submersion is the Hopf map  $H : S^3 \subset \mathbb{C}^2 \rightarrow S^2(\frac{1}{2}) \subset \mathbb{R} \times \mathbb{C}$ :

$$H(z, w) = \left(\frac{1}{2}(|z|^2 - |w|^2), z\bar{w}\right). \quad (24)$$

It follows:

$$H\left(\frac{1}{\sqrt{2}}\mathcal{F}(t)\right) = \left(0, \frac{1}{2}T(t)\bar{N}(t)\right) = \left(0, -\frac{i}{2}\right) = H\left(\frac{1}{\sqrt{2}}\mathcal{E}(t)\right). \quad (25)$$

Hence, considering  $H$  as a projection map of the  $S^1$ -principal bundle  $S^3 \rightarrow S^2(\frac{1}{2})$  we have that  $\frac{1}{\sqrt{2}}\mathcal{F}$  and  $\frac{1}{\sqrt{2}}\mathcal{E}$  belong to the same fiber, namely that over the South pole of the sphere  $S^2(\frac{1}{2})$ .

iv) Suppose now that our curve  $C$  belongs to the plane  $xOz$  of the physical space  $\mathbb{R}^3$  as  $C : r(t) = (f(t), 0, F(t))$  with  $f > 0$  on  $I$  and consider the rotational surface generated by  $C$  as:

$$\Sigma : \bar{r}(t, \varphi) := (f(t) \cos \varphi, f(t) \sin \varphi, F(t)), \quad \varphi \in S^1.$$

Its principal curvatures depend only on  $t$ , [7, p. 85]:

$$k_1 = k, \quad k_2 = \frac{F'}{\|r'\|f} \quad (26)$$

and then for  $F' = f$  we have that  $k_f$  of  $C$  is exactly the difference  $k_1 - k_2$  of the principal curvatures of  $\Sigma$ ; consequently the umbilic circles of  $\Sigma$  are provided by the zeros of  $k_f$  and are parametrized by  $\varphi \in S^1$ .

For  $F' = f$  the curvatures of  $C$  are expressed only through the function  $F$  as:

$$k(t) = \frac{[F''(t)]^2 - F'(t)F'''(t)}{[F'(t)^2 + F''(t)^2]^{\frac{3}{2}}}, \quad k_f(t) = \frac{-F'(t)F'''(t) - [F'(t)]^2}{[F'(t)^2 + F''(t)^2]^{\frac{3}{2}}} \quad (27)$$

and due to the presence of the third derivative of  $F$  we recall its Schwarzian derivative:

$$S_F = \frac{F'''}{F'} - \frac{3}{2} \left( \frac{F''}{F'} \right)^2 \quad (28)$$

which implies the new formulae:

$$k = \frac{(F'')^2 - 2(F')^2 S_F}{2[(F')^2 + (F'')^2]^{\frac{3}{2}}}, \quad k_f = \frac{-3(F'')^2 - 2(F')^2 S_F - 2(F')^2}{2[(F')^2 + (F'')^2]^{\frac{3}{2}}}. \quad (29)$$

In conclusion, a smooth  $F$  with negative Schwarzian derivative will give a positive curvature  $k$  for  $C$  while a positive Schwarzian derivative  $S_F$  produces a negative flow-curvature  $k_f$ .

v) The nature and the relationship between our frames can be put in the framework of moving frames of [7, p. 32]. Recall that the set of all orientation-preserving Euclidean isometries forms a Lie group,  $E(2) := \mathbb{R}^2 \times SO(2)$ , with the standard projection  $\pi_1$  on the first factor making  $E(2) \rightarrow \mathbb{R}^2$  an  $S^1$ -principal bundle. A moving frame along  $C$  is a map  $F : I \rightarrow E(2)$  such that  $\pi_1 \circ F = r$ . But  $C$  defines also a 1-parameter family of bijections of  $SO(2)$ :

$$L^C : I \rightarrow \text{Bijections}(SO(2)), t \rightarrow L^C(t) : SO(2) \rightarrow SO(2), A \rightarrow R(t)A, (L^C(t))^{-1} = L^C(-t).$$

Then our frames are  $\mathcal{F} : I \rightarrow E(2)$  as  $\mathcal{F}(t) = (r(t), T(t), N(t))$  and  $\mathcal{E} : I \rightarrow E(2)$  as  $\mathcal{E}(t) = (r(t), (L^C(t) \circ \pi_2 \circ \mathcal{F})(t))$ .

vi) Suppose now that the curve  $C$  is in the space  $\mathbb{R}^3$  and is bi-regular; hence it has the Frenet frame  $(T, N, B)$  and the pair (curvature, torsion) =  $(k, \tau)$ . We define its flow-frame as:

$$\begin{pmatrix} T \\ E_2^f \\ E_3^f \end{pmatrix} (t) := \begin{pmatrix} 1 & 0_2(h) \\ 0_2(v) & R(t) \end{pmatrix} \begin{pmatrix} T \\ N \\ B \end{pmatrix}, \quad 0_2(h) := (0, 0), \quad 0_2(v) := \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

and then, its matrix moving equation is:

$$\frac{d}{dt} \begin{pmatrix} T \\ E_2^f \\ E_3^f \end{pmatrix} (t) = \|r'(t)\| \begin{pmatrix} 0 & k_f^2(t) & k_f^3(t) \\ -k_f^2(t) & 0 & \tau_f(t) \\ -k_f^3(t) & -\tau_f(t) & 0 \end{pmatrix} \begin{pmatrix} T \\ E_2^f \\ E_3^f \end{pmatrix} (t).$$

A similar computation yields:

$$k_f^2(t) = k(t) \cos t, \quad k_f^3(t) = k(t) \sin t, \quad \tau_f(t) = \tau(t) - \frac{1}{\|r'(t)\|} < \tau(t).$$

We point out the formal similarity with the Darboux equations of a curve on a given surface and then a curve  $C$  with vanishing  $\tau_f$  will be called flow-geodesic in  $\mathbb{R}^3$ . Hence, if  $C$  is naturally parametrized then  $C$  is a flow-geodesic if and only if its torsion has the constant value 1; for this class of space curves and examples see [1]. In order to express the above moving equation in the compact form as in the theory of space curves:

$$\omega_f(t) \times T(t) = T'(t), \quad \omega_f(t) \times E_2^f(t) = (E_2^f)'(t), \quad \omega_f(t) \times E_3^f(t) = (E_3^f)'(t)$$

we associate a vector field along  $C$ , called flow-Darboux:

$$\omega_f(t) := \|\gamma'(t)\|[\tau_f(t)T(t) - k_f^3(t)E_2^f(t) + k_f^2(t)E_3^f(t)].$$

Something similar but with the rotation with respect to an angle  $\theta = \theta(s)$  appears in [12] under the name of quasi frame for  $C$ . Our choice corresponds to the angle  $\theta(s) = -s$ .

vii) Suppose that the curvature function  $t \rightarrow k(t)$  is always strictly positive (or strictly negative). Then the evolute of  $C$  is the curve:

$$C_e : r_e(t) := r(t) + \frac{1}{k(t)}N(t).$$

With this model in mind, for a non-flat-flow curve we associate its flow-evolute as being the curve:

$$C_{fe} : r_{fe}(t) := r(t) + \frac{1}{k_f(t)}E_2^f(t).$$

We will obtain this curve for some examples below. So, the line  $C$  discussed in the example 1i has the flow-evolute

$$C_{fe} : r_{fe}(t) = r_0 + (t - \sin t)u - \cos t(iu)$$

and for  $r_0 = (0, 1) = iu$  this last curve is exactly the cycloid of radius  $R = 1$  according to the example 3 below.  $\square$

Returning to the plane curves let  $J \subseteq \mathbb{R}$  be another open interval and fix the diffeomorphism  $\varphi : s \in J \rightarrow t \in I$  with the smooth inverse  $\varphi^{-1} : t \in I \rightarrow s \in J$ . Since  $r'(s) = \varphi'(s)r'(t(s))$  we restrict our study to the class  $Diff_+(J, I)$  of orientation-preserving diffeomorphisms:  $\varphi'(s) > 0$ , for all  $s \in J$ . The transformation of the flow-curvature under the action of  $\varphi$  is:

$$k_f(s) = k(t) - \frac{1}{\varphi'(s)\|r'(t)\|} \quad (30)$$

and then:

$$k_f(s) - k_f(t) = \frac{1}{\|r'(t)\|} \left[ 1 - \frac{1}{\varphi'(s)} \right]. \quad (31)$$

**Proposition 2.** (the rigidity of the flow-curvature) *The only orientation-preserving diffeomorphism  $\varphi$  which preserves also the flow-curvature of  $C$  is an interval shift on the real line  $\varphi(s) = s + s_0$ ,  $s_0 \in (0, +\infty)$ .*

A natural important problem is the class of curves with prescribed flow-curvature. For example, if we ask the vanishing of the flow-curvature for a graphic curve  $C_F : r(t) = (t, F(t))$  then it follows the differential equation:

$$\frac{F''(t)}{[1 + (F'(t))^2]^{\frac{3}{2}}} = \frac{1}{[1 + (F'(t))^2]^{\frac{1}{2}}}. \quad (32)$$

Since this equation reads:

$$\frac{F''(t)}{1 + (F'(t))^2} = 1 \quad (33)$$

we have exactly the Grim Reaper solution, [3, p. 28], a famous solution of the curve shortening flow:

$$F_u(t) = u - \ln(\cos t), \quad t \in \left(-\frac{\pi}{2}, \frac{\pi}{2}\right), \quad u \in \mathbb{R} \quad (34)$$

with the usual curvature  $k(t) = \cos t$  and the frames:

$$\mathcal{F}(t) = \begin{pmatrix} e^{it} \\ e^{i(t+\frac{\pi}{2})} \end{pmatrix}, \quad \mathcal{E} = \begin{pmatrix} (1, 0) = \bar{i} \\ (0, 1) = \bar{j} \end{pmatrix} = \text{constant}. \quad (35)$$

Another formalism is that of [14, p. 2] if  $r : S^1 \simeq [0, 2\pi) \rightarrow \mathbb{R}^2$  is naturally parametrized then there exists the smooth function  $\theta : S^1 \rightarrow \mathbb{R}$ , called *normal angle*, such that:

$$N(s) = e^{i\theta(s)} = (\cos \theta(s), \sin \theta(s)), \quad T(s) = -iN(s) = -ie^{i\theta(s)} = e^{i(\theta(s) - \frac{\pi}{2})} \quad (36)$$

and then the Frenet equations yield:

$$\frac{d\theta}{ds}(s) = k(s). \quad (37)$$

In conclusion, the constant value  $\beta \in \mathbb{R}$  of the flow-curvature of a closed convex curve means  $\theta(s) = (\beta + 1)s + \alpha$  for all  $s \in S^1$  with  $\alpha \in \mathbb{R}$  an arbitrary constant. The flow-frame corresponding to the equations (36) is:

$$E_1^f(s) = (\sin(\theta(s) - t(s)), -\cos(\theta(s) - t(s))), \quad E_2^f(s) = (\cos(\theta(s) - t(s)), \sin(\theta(s) - t(s))) \quad (38)$$

which, in turn, is the Frenet frame of a new curve with the same natural parameter  $s$  but having the normal angle  $\tilde{\theta}(s) := \theta(s) - t(s)$ .

The formula (37) can be replaced with  $\frac{d(\theta - \pi/2)}{ds}(s) = k(s)$  which expresses the curvature  $k$  as the derivative of the angle between  $T \in \mathcal{X}_C$  and the unit vector  $\bar{i}$ . Following this approach the paper [6] generalizes  $k$  to a curvature-type function  $k_V$  defined with respect to an arbitrary  $V \in \mathcal{X}_C$ . A main result of the cited work is that  $k_V = k_W$  if and only if the angle between  $V$  and  $W$  is constant along  $C$ . Hence, we can apply the last statement of the Remark ii) and then two flow-Fermi-Walker unit vectors  $V, W \in \mathcal{X}_C$  yield the same curvature-type function.

In the following we present a couple of examples in order to remark the computational aspects of our approach.



**Example 2.** The involute of the unit circle  $S^1$  is:

$$C : r(t) = (\cos t + t \sin t, \sin t - t \cos t) = (1 - it)e^{it}, \quad t \in (0, +\infty). \quad (39)$$

A direct computation gives:

$$r'(t) = (t \cos t, t \sin t) = te^{it}, \quad \|r'(t)\| = t, \quad k(t) = \frac{1}{t} > 0, \quad (40)$$

and then this curve is also a flat-flow one and having the same flow-frame as the Grim Reaper. This example can be treated also with respect to a natural parameter  $s \in (0, +\infty)$  which is provided by  $t := \sqrt{2s}$ . For example, the normal angle function is  $\theta(s) = \frac{\pi}{2} + \sqrt{2s}$  since then  $r'(s) = e^{i\sqrt{2s}}$ . Comparing with the approach above it results the constants  $\alpha = \frac{\pi}{2}$  and  $\beta = \sqrt{2} - 1$ .  $\square$

**Example 3.** Recall that for  $R > 0$  the cycloid of radius  $R$  has the equation:

$$C : r(t) = R(t - \sin t, 1 - \cos t) = R[(t, 1) - e^{i(\frac{\pi}{2}-t)}], \quad t \in \mathbb{R}. \quad (41)$$

Remark that here we have a twisted situation of the Remark iv) namely the derivative of the first component of the vector  $r(t)$  is exactly the second component. The Schwarzian derivative is:

$$S_{t-\sin t}(t) = \frac{\cos t}{\sin t} - \frac{3}{2} \left( \frac{\cos \frac{t}{2}}{\sin \frac{t}{2}} \right)^2, \quad t \in \mathbb{R} \setminus \mathbb{Z}\pi. \quad (42)$$

We have immediately:

$$r'(t) = R(1 - \cos t, \sin t) = R[(1, 0) - e^{it}], \quad \|r'(t)\| = 2R|\sin \frac{t}{2}|, \quad k(t) = -\frac{1}{4R|\sin \frac{t}{2}|}, \quad (43)$$

and then we restrict our definition domain to  $(0, \pi)$ . It follows:

$$\begin{cases} k_f(t) = -\frac{3}{4R \sin \frac{t}{2}} < 0, \\ E_1^f(t) = (\sin \frac{3t}{2}, \cos \frac{3t}{2}) = e^{i(\frac{\pi}{2}-\frac{3t}{2})}, E_2^f(t) = (-\cos \frac{3t}{2}, \sin \frac{3t}{2}) = e^{i(\pi-\frac{3t}{2})}. \end{cases} \quad (44)$$

Again a natural parameter  $s$  is provided by:  $t = 2 \arccos(1 - \frac{s}{4R})$  and the flow-evolute of  $C$  is the curve:

$$C_{fe} : r_{fe}(t) = R(t - \sin t, 1 - \cos t) + \frac{4}{3}R \sin \frac{t}{2}(\cos t, -\sin t), \quad t \in (0, \pi).$$

$\square$

**Example 4.** The derivative curve  $r'$  from (40) is an Archimedes' spiral. This spiral is given in polar coordinates as:

$$A(\text{spiral}) : \rho(t) = Rt, \quad R > 0 \quad (45)$$

and hence:

$$k_f(t) = \frac{1}{R(t^2 + 1)^{\frac{3}{2}}} > 0 \quad (46)$$

while its flow-evolute is the curve:

$$C_{fe} : r_{fe}(t) = R(t \cos t, t \sin t) + R(1 + t^2)(-\sin t - t \cos t, \cos t - t \sin t).$$

□

**Example 5.** Fix  $\alpha \in \mathbb{R}^*$  and the naturally parametrized curve  $C$ . Then the  $\alpha$ -parallel curve of  $C$  is the new curve:

$$C_\alpha : \tilde{r}(t) := r(t) + \alpha N(t), \quad t \in I \quad (47)$$

with:

$$\tilde{T}(t) = \frac{1 - \alpha k(t)}{|1 - \alpha k(t)|} T(t), \quad \tilde{N}(t) = \frac{1 - \alpha k(t)}{|1 - \alpha k(t)|} N(t), \quad \tilde{k}(t) = k(t). \quad (48)$$

Hence, we consider that  $\alpha$  does not belongs to the range of the function  $\frac{1}{k}$  and the new flow-curvature is:

$$\tilde{k}_f(t) = k(t) - \frac{1}{|1 - \alpha k(t)|}. \quad (49)$$

□

We finish this note with the problem raised in the beginning, namely the possible variants of the curve shortening flow. Recall that the setting of this question consists in a 1-parameter family of plane curves  $C_u : r = r_u(t) = r(t, u)$  satisfying:

$$\frac{\partial r(t, u)}{\partial u} = k(t, u)N(t, u). \quad (50)$$

It follows immediately an expression in terms of flow-apparatus:

$$\frac{\partial r(t, u)}{\partial u} = \left( k_f(t, u) + \frac{1}{\|r'(t, u)\|} \right) [-\sin t E_1^f(t, u) + \cos t E_2^f(t, u)]. \quad (51)$$

The first variant which we propose as an open problem is to study the flow-variant of (50):

$$\frac{\partial r(t, u)}{\partial u} = k_f(t, u)E_2^f(t, u). \quad (52)$$

The second variant is to generalize all this study through a general smooth function  $\Omega \in C^\infty(\mathbb{R})$ . More precisely, we use the equation (10) with  $R$  replaced by  $R \circ \Omega$  to define the notion of  $\Omega$ -frame for the plane curve  $C$ ; we note that for a particular choice of  $\Omega$  the 3-dimensional variant of the remark vi) is called *positional adapted frame* in [11]. Then the  $\Omega$ -curvature of the plane curve  $C$  is:

$$k_\Omega(t) = k(t) - \frac{\Omega'(t)}{\|r'(t)\|} \quad (53)$$

and the curves in polar coordinates with vanishing  $\Omega$ -curvature are provided by:

$$\rho(t) = R e^{\int_{t_0}^t \cot[\Omega(u) - u + C] du}, \quad R > 0, \quad C \in \mathbb{R}. \quad (54)$$

The flow-curvature corresponds to the identity map  $\Omega = 1_{\mathbb{R}}$ . Moreover, if  $C$  is naturally parametrized then  $k_{\Omega} = (\theta - \Omega)'$  which means that the case  $\Omega = \theta + \text{constant}$  provides a zero  $\Omega$ -curvature.

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#### REFERENCES

- [1] Bates, L. M., Melko, O. M., On curves of constant torsion I, *J. Geom.*, 104 (2) (2013), 213–227. <https://doi.org/10.1007/s00022-013-0166-2>
- [2] Bishop, R. L., There is more than one way to frame a curve, *Am. Math. Mon.*, 82 (1975), 246–251. <https://doi.org/10.2307/2319846>
- [3] Chou, K.-S., Zhu, X.-P., The curve shortening problem, Boca Raton, FL: Chapman & Hall/CRC, 2001. Zbl 1061.53045
- [4] Crasmareanu, M., The flow-curvature of spacelike parametrized curves in the Lorentz plane, *Proceedings of the International Geometry Center*, 15 (2) (2022), 100–108. <https://doi.org/10.15673/tmgc.v15i2.2281>
- [5] Crasmareanu, M., Frigioiu, C. Unitary vector fields are Fermi-Walker transported along Rytov-Legendre curves, *Int. J. Geom. Methods Mod. Phys.*, 12 (10) (2015), , Article ID 1550111. <https://doi.org/10.1142/S021988781550111X>
- [6] Gózdź, S., Curvature type functions for plane curves, *An. Științ. Univ. Al. I. Cuza Iași Mat.*, 39 (3) (1993), 295–303. Zbl 0851.53001
- [7] Jensen, G. R., Musso, E., Nicolodi, L., Surfaces in classical geometries. A treatment by moving frames, Universitext, Springer, 2016. Zbl 1347.53001
- [8] Mazur, B., Perturbations, deformations, and variations (and “near-misses”) in geometry, physics, and number theory, *Bull. Am. Math. Soc.*, 41 (3) (2004), 307–336. <https://doi.org/10.1090/S0273-0979-04-01024-9>
- [9] Miron, R., Une généralisation de la notion de courbure de parallélisme, *Gaz. Mat. Fiz., București, Ser. A* 10 (63) (1958), 705–708. Zbl 0087.36101
- [10] Miron, R., The geometry of Myller configurations. Applications to theory of surfaces and nonholonomic manifolds, Bucharest: Editura Academiei Române, 2010. Zbl 1206.53003
- [11] Özen, K. E., Tosun, M., A new moving frame for trajectories with non-vanishing angular momentum, *J. of Mathematical Sciences and Modelling*, 4 (1) (2021), 7–18. <https://doi.org/10.33187/jmsm.869698>
- [12] Soliman, M. A., Nassar, H.A.-A., Hussien, R. A., Youssef, T., Evolutions of the ruled surfaces via the evolution of their directrix using quasi frame along a space curve, *J. of Applied Mathematics and Physics*, 6 (2018), 1748–1756. <https://doi.org/10.4236/jamp.2018.68149>
- [13] Younes, L., Shapes and diffeomorphisms. 2nd updated edition, Applied Mathematical Sciences 171. Berlin: Springer, 2019. Zbl 1423.53002
- [14] Zhu, X.-P., Lectures on mean curvature flows, AMS/IP Studies in Advanced Mathematics vol. 32, Providence, RI: American Mathematical Society, 2002. Zbl 1197.53087