

DIRAC STRUCTURES FROM LIE INTEGRABILITY

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We prove that a pair ($F =$ vector sub-bundle of TM , its annihilator) yields an almost Dirac structure which is Dirac if and only if F is Lie integrable. Then a flat Ehresmann connection on a fiber bundle ξ yields two complementary, but not orthogonally, Dirac structures on the total space M of ξ . These Dirac structures are also Lagrangian sub-bundles with respect to the natural almost symplectic structure of the big tangent bundle of M . The tangent bundle in Riemannian geometry is discussed as particular case and the 3-dimensional Heisenberg space is illustrated as example. More generally, we study the Bianchi–Cartan–Vranceanu metrics and their Hopf bundles.

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Fix M a smooth and connected $m \geq 2$ -dimensional manifold with the tangent bundle TM and cotangent bundle T^*M . The $C^\infty(M)$ -module of sections of these bundles are: $\mathcal{X}(M) = \Gamma(TM)$ the Lie algebra of vector fields on M and $\Omega^1(M) = \Gamma(T^*M)$ the real linear space of 1-forms on M .

The manifold $T^{\text{big}}M := TM \oplus T^*M$ is called *big tangent bundle* in [8] and *the Pontryagin bundle* of M in [5]; for simplicity we will use the first name. This manifold is the total space of a vector bundle $\pi : T^{\text{big}}M \rightarrow TM$ with π the first projection and $\Gamma(T^{\text{big}}M) = \mathcal{X}(M) \otimes \Omega^1(M)$ is endowed with *the Courant structure* $(\langle, \rangle, [,])$, [2]:

(1) the (neutral) inner product:

$$\langle (X, \alpha), (Y, \beta) \rangle = \frac{1}{2}(\beta(X) + \alpha(Y)), \quad (1)$$

(2) the (skew-symmetric) Courant bracket:

$$[(X, \alpha), (Y, \beta)] = \left([X, Y], \mathcal{L}_X \beta - \mathcal{L}_Y \alpha - \frac{1}{2} d(\beta(X) - \alpha(Y)) \right), \quad (2)$$

where \mathcal{L}_X is the Lie derivative with respect to the vector field $X \in \mathcal{X}(M)$.

Definition 1 ([2]). A vector sub-bundle L of $T^{\text{big}}M$ which is \langle, \rangle -isotropic and of maximal rank m is called *almost Dirac structure* on M . If, in addition, the $C^\infty(M)$ -module $\Gamma(L)$ of sections of L is $[\cdot, \cdot]$ -integrable then L is a *Dirac structure*.

Two associated notions are:

- (1) a pair of complementary Dirac subspaces is called a *reflector* ([4, 8]) or *Manin triple* in [6],
- (2) if L is a Dirac structure on M then the triple $(L, [\cdot, \cdot]_L, pr : L \rightarrow TM)$ is a *Lie algebroid* on M ([2, Theorem 2.3.4]).

Fix F a vector sub-bundle of TM and denote by F^0 its annihilator, that means the vector sub-bundle of T^*M which vanish on F . Then $L_F = F \oplus F^0$ is a vector sub-bundle of rank m of $T^{\text{big}}M$.

The main result of this short note is a version of [8, Example 2.3].

Theorem 2. L_F is an almost Dirac structure which is Dirac on M if and only if F is Lie integrable.

Proof. L_F is maximally isotropic with respect to the inner product (1). Suppose now that F is Lie integrable; then it defines a foliation \mathcal{F} of M . The elements of $\Gamma(F^0)$ are 1-forms which vanish by restriction to each leaf of \mathcal{F} . Also, their exterior derivative induce zero 2-forms on the leaves of \mathcal{F} . Using these facts, it results that the Courant bracket of any two sections of L_F is:

$$[(X, \alpha), (Y, \beta)] = ([X, Y], 0),$$

which yields the conclusion about the Dirac structure of L_F .

Conversely, if L_F is a Dirac structure from the Lie algebroid structure of L_F we get that $pr(L_F)$ is an integrable generalized distribution in the Stefan–Sussman sense. But here $pr(L_F)$ being exactly F is a true vector sub-bundle of TM i.e. of constant rank and then the Stefan–Sussman integrability is exactly the Lie-Frobenius one. \square

Example 3 (Ehresmann connection). Let $\xi = (M, \pi, B, F)$ be a fiber bundle; we use the approach of [1, p. 105]. Let n be the dimension of B and $x = (x^i; 1 \leq i \leq n)$, respectively $y = (y^\alpha; 1 \leq \alpha \leq m - n)$ the coordinates on B respectively F . Then, $(x, y) = (x^i, y^\alpha)$ are the coordinates on M . Applying the tangent functor to the bundle ξ we get $T\pi : TM \rightarrow TB$ and then the kernel of $T\pi$ is a vector bundle over M denoted $T_{V,\pi}M$ and called *the vertical bundle of ξ* . Its space of sections $\Gamma(T_{V,\pi}M)$ is spanned by $\left\{ \frac{\partial}{\partial y^\alpha} \right\}$.

Let an *Ehresmann connection* E on ξ with the horizontal vector sub-bundle $T_{H,\pi}M$ over M , complementary to $T_{V,\pi}M$ in TM , [1, p. 107]. It results the existence of smooth functions $\{N_i^\alpha(x, y)\}$ on M , called *the coefficients of connection*, such that $\Gamma(T_{H,\pi}M)$ is spanned by $\{\frac{\delta}{\delta x^i} = \frac{\partial}{\partial x^i} - N_i^\alpha \frac{\partial}{\partial y^\alpha}\}$.

Let $T_{V,\pi}^*M$ be the annihilator of $F_1 = T_{V,\pi}M$; then $\Gamma(T_{V,\pi}^*M)$ is spanned by $\{dx^i\}$. Similarly, let $T_{H,\pi}^*M$ be the annihilator of $F_2 = T_{H,\pi}M$; it results $\Gamma(T_{H,\pi}^*M)$ is spanned by $\{\delta y^\alpha = dy^\alpha + N_i^\alpha dx^i\}$. Our main objects in the following are the vector bundles over $M: L_{F_1} = V_D(\pi, E) = (T_{V,\pi}M, T_{V,\pi}^*M)$, respectively $L_{F_2} = H_D(\pi, E) = (T_{H,\pi}M, T_{H,\pi}^*M)$. The rank of both is m .

We compute:

$$\mathcal{L}_{\frac{\delta}{\delta x^i}} \delta y^\alpha = R_{ik}^\alpha dx^k - \frac{\partial N_i^\alpha}{\partial y^\rho} \delta y^\rho, \quad \mathcal{L}_{\frac{\partial}{\partial y^\alpha}} dx^j = 0, \quad (3)$$

where R is *the curvature* of the Ehresmann connection E :

$$R_{ij}^\alpha = \frac{\delta N_i^\alpha}{\delta x^j} - \frac{\delta N_j^\alpha}{\delta x^i} \quad (4)$$

provided by the condition of Lie integrability of the distribution $\Gamma(T_{H,\pi}M)$:

$$\left[\frac{\delta}{\delta x^i}, \frac{\delta}{\delta x^j} \right] = R_{ij}^\alpha \frac{\partial}{\partial y^\alpha}. \quad (5)$$

Since $T_{V,\pi}M$ is integrable it results:

Corollary 4. $V_D(\pi, E)$ is a Dirac structure on M . $H_D(\pi, E)$ is an almost Dirac structure which is Dirac if and only if the connection E is flat.

It results that a flat connection yields a Manin triple:

$$T^{\text{big}}M = V_D(\pi, E) \oplus H_D(\pi, E) \quad (6)$$

in analogy with the decomposition of the tangent bundle provided by the Ehresmann connection:

$$TM = T_{V,\pi}M \oplus T_{H,\pi}M. \quad (7)$$

From:

$$\left\langle \left(\frac{\partial}{\partial y^\alpha}, dx^i \right), \left(\frac{\delta}{\delta x^j}, \delta y^\beta \right) \right\rangle = \frac{1}{2}(\delta_j^i + \delta_\alpha^\beta) \quad (8)$$

we get that the decomposition (6) is not \langle, \rangle -orthogonally. Therefore, we obtain two Lie algebroids over M provided by a flat connection.

On this example we will treat other two main constructions of [2]. Firstly, let $\pi^* : T^{\text{big}}TM \rightarrow T^*M$ be the canonical second projection. Relation (2.2.3 – 4) of [2, p. 645] states that for a general almost Dirac structure we have: $\pi(L) = L \cap T^*M$ and $\pi^*(L) = (L \cap TM)^0$ and a skew-symmetric map $\Omega : \pi(L) \rightarrow \pi^*(L)$ is derived with the formula: $\Omega(\pi(u)) = \pi^*(u)|_{\pi(L)}$. In our framework we get: $\pi(L_F) = F$ and

$\pi^*(L_F) = F^0$. Also, we have:

- (1) $\Omega_V : T_{V,\pi}M \rightarrow T_{V,\pi}^*M$ is $\Omega(\frac{\partial}{\partial y^\alpha}|_{(x,y)}) = dx^i|_{(x,y)}$,
- (2) $\Omega_H : T_{H,\pi}M \rightarrow T_{H,\pi}^*M$ is $\Omega(\frac{\delta}{\delta x^i}|_{(x,y)}) = \delta y^\alpha|_{(x,y)}$,

for every $(x, y) \in M$ since as it is pointed out several times in [2] in the manifold case we need to work pointwise.

Secondly, a function $f \in C^\infty(M)$ is called *admissible* ([2, p. 650]) or *Hamiltonian* ([8]) if there exists a vector field $X_f \in \mathcal{X}(M)$ such that $(X_f, df) \in \Gamma(L)$ and then X_f is called *the Hamiltonian vector field* of f and a bracket on admissible functions is given by $\{f, g\} = X_f(g)$. So, for $T_{V,\pi}M$ we get that the coordinate functions x^i are admissible functions and the bracket is trivial:

$$\{x^i, x^j\} = \frac{\partial x^j}{\partial y^\alpha} = 0, \quad (9)$$

where $\frac{\partial}{\partial y^\alpha}$ is the Hamiltonian vector field of $f = x^i$ since $(\frac{\partial}{\partial y^\alpha}, dx^i) \in \Gamma(T_{V,\pi}M)$. $T_{H,\pi}M$ has $\{y^\alpha\}$ as admissible functions only if the Ehresmann connection has vanishing coefficients, i.e. $N_i^\alpha = 0$, in which case $\frac{\delta}{\delta x^i} = \frac{\partial}{\partial x^i}$ are the corresponding Hamiltonian vector fields and again the bracket is trivial.

Particular Case 5 (Riemannian geometry). Suppose that ξ is exactly the tangent bundle of B i.e. $M = TB$. Then $m = 2n$, $F = \mathbb{R}^n$ and $\alpha = k$. Every Riemannian metric g on B gives rise to the Levi-Civita connection E where $N_i^k = \Gamma_{ia}^k(x)y^a$, with (Γ_{ij}^k) being the Christoffel symbols of g . Then, a flat Riemannian metric on B yields two complementary Dirac structures on TB .

Let us remark that some classes of Dirac structures naturally associated to Lagrangian systems appear in [11].

Remark 6 (The almost symplectic point of view). Let us recall that the big tangent bundle has also a non-degenerate skew-symmetric 2-form, [8]:

$$\Omega((X, \alpha), (Y, \beta)) = \frac{1}{2}(\beta(X) - \alpha(Y)). \quad (10)$$

From:

$$\Omega\left(\left(\frac{\partial}{\partial y^\alpha}, dx^i\right), \left(\frac{\partial}{\partial y^\beta}, dx^j\right)\right) = 0, \quad \Omega\left(\left(\frac{\delta}{\delta x^i}, \delta y^\alpha\right), \left(\frac{\delta}{\delta x^j}, \delta y^\beta\right)\right) = 0 \quad (11)$$

we get that the complementary Dirac structures $V_D(\pi, E)$ and $H_D(\pi, E)$ are Ω -Lagrangian sub-bundles of $T^{\text{big}}M$.

Example 7. Let H_3 be the 3-dimensional *Heisenberg space*, [3]; recall that $H_3 = (\mathbb{R}^3, g)$ with the Riemannian metric:

$$g = dx^2 + dy^2 + \left(dz + \frac{y}{2}dx - \frac{x}{2}dy\right)^2. \quad (12)$$

An orthonormal basis of left-invariant vector fields on this Lie group is:

$$X_1 = \frac{\partial}{\partial x} - \frac{y}{2} \frac{\partial}{\partial z}, \quad X_2 = \frac{\partial}{\partial y} + \frac{x}{2} \frac{\partial}{\partial z}, \quad X_3 = \frac{\partial}{\partial z} \quad (13)$$

with the dual frame:

$$\theta_1 = dx, \quad \theta_2 = dy, \quad \theta_3 = dz + \frac{y}{2}dx - \frac{x}{2}dy. \quad (14)$$

Let (x, y, z, X, Y, Z) be the coordinates on TH_3 and (x, y, z, p_x, p_y, p_z) the coordinates on T^*H_3 ; then $T^{\text{big}}H_3$ has the coordinates $(x, y, z, X, Y, Z, p_x, p_y, p_z)$. Let the locally trivial fiber bundle $\pi : M = H_3 \rightarrow B = \mathbb{R}^2$, $\pi(x, y, z) = (x, y)$, then $m = 3$ and $n = 2$. Fix the connection $E = E_g$ such that $T_{H,\pi}H_3$ is the g -orthogonal complement of $T_{V,\pi}H_3$.

Then $\Gamma(T_{V,\pi}H_3)$ is spanned by X_3 while $\Gamma(T_{V,\pi}^*H_3)$ is spanned by $\{\theta_1, \theta_2\}$; $\Gamma(T_{H,\pi}H_3)$ is spanned by $\{\frac{\delta}{\delta x} = X_1, \frac{\delta}{\delta y} = X_2\}$ and $\Gamma(T_{H,\pi}^*H_3)$ is spanned by $\{\theta_3\}$. We get the coefficients of the connection:

$$N_1^1 = \frac{y}{2}, \quad N_2^1 = -\frac{x}{2} \quad (15)$$

and therefore the curvature is:

$$R_{12}^1 = -R_{21}^1 = 1. \quad (16)$$

(Of course, we know that g is not a flat metric but we include these computations in order to illustrate completely this example.) In conclusion, the only Dirac structure provided by Corollary 4 on H_3 is:

$$V_D(\pi, E_g) = \{(x, y, z, 0, 0, Z, p_x, p_y, 0) \in T^{\text{big}}H_3\} \quad (17)$$

while:

$$H_D(\pi, E_g) = \left\{ \left(x, y, z, X - \frac{y}{2}Z, Y + \frac{x}{2}Z, 0, 0, 0, p_z + \frac{y}{2}p_x - \frac{x}{2}p_y \right) \in T^{\text{big}}H_3 \right\} \quad (18)$$

is an almost Dirac structure on H_3 .

Example 8. Fix k and τ two real numbers and denotes by M_k^3 the manifold $\{(x, y, z) \in \mathbb{R}^3; F(x, y, z) = 1 + k(x^2 + y^2) > 0\}$. We shall consider on M_k^3 the *Bianchi-Cartan-Vranceanu metric*, [9, p. 343]:

$$g_{k,\tau} = \frac{1}{F^2}dx^2 + \frac{1}{F^2}dy^2 + \left(dz + \frac{\tau y}{2F}dx - \frac{\tau x}{2F}dy \right)^2. \quad (19)$$

For $k = 0$ and $\tau = 1$ we get the previous example; for other remarkable examples as well as the history of these metrics see [9].

An orthonormal basis in $\mathcal{X}(M_k^3)$ is:

$$X_1 = F \frac{\partial}{\partial x} - \frac{\tau y}{2} \frac{\partial}{\partial z}, \quad X_2 = F \frac{\partial}{\partial y} + \frac{\tau x}{2} \frac{\partial}{\partial z}, \quad X_3 = \frac{\partial}{\partial z} \quad (20)$$

with the dual frame:

$$\theta_1 = \frac{dx}{F}, \quad \theta_2 = \frac{dy}{F}, \quad \theta_3 = dz + \frac{\tau y}{2F}dx - \frac{\tau x}{2F}dy. \quad (21)$$

Let $M_k^2 = \{(x, y) \in \mathbb{R}^2; F(x, y) > 0\}$ with the metric $g_k = \frac{dx^2 + dy^2}{F^2}$. Then the map $\pi : (M_k^3, g_{k,\tau}) \rightarrow (M_k^2, g_k)$, $\pi(x, y, z) = (x, y)$ being a Riemannian submersion yields

a fiber bundle called *Hopf bundle* in [9, p. 344] since for $k = \frac{\tau^2}{4}$ we get the Hopf fibration $\pi : S^3(k) \rightarrow S^2(4k)$. Again, let the connection $E = E_g$ such that $T_{H,\pi}M_k^3$ is the $g_{k,\tau}$ -orthogonal complement of $T_{V,\pi}M_k^3$.

Then $\Gamma(T_{V,\pi}M_k^3)$ is spanned by X_3 while $\Gamma(T_{V,\pi}^*M_k^3)$ is spanned by $\{\delta x = F\theta_1, \delta y = F\theta_2\}$; $\Gamma(T_{H,\pi}M_k^3)$ is spanned by $\{\frac{\delta}{\delta x} = \frac{1}{F}X_1, \frac{\delta}{\delta y} = \frac{1}{F}X_2\}$ and $\Gamma(T_{H,\pi}^*M_k^3)$ is spanned by $\{\theta_3\}$. We get the coefficients of the connection:

$$N_1^1 = \frac{\tau y}{2F}, \quad N_2^1 = -\frac{\tau x}{2F} \tag{22}$$

and therefore the curvature is:

$$R_{12}^1 = -R_{21}^1 = \frac{\tau}{F^2}. \tag{23}$$

In conclusion, the only Dirac structure provided by Corollary 4 on M_k^3 is:

$$V_D(\pi, E_g) = \{(x, y, z, 0, 0, Z, p_x, p_y, 0) \in T^{\text{big}}M_k^3\} \tag{24}$$

while:

$$H_D(\pi, E_g) = \left\{ \left(x, y, z, X - \frac{\tau y}{2F}Z, Y + \frac{\tau x}{2F}Z, 0, 0, 0, p_z + \frac{\tau y}{2F}p_x - \frac{\tau x}{2F}p_y \right) \in T^{\text{big}}M_k^3 \right\} \tag{25}$$

is an almost Dirac structure on M_k^3 .

Remark 9 (A relationship with previous works). Returning to the general framework, let us recall that a connection on ξ is equivalent with an almost product structure P on M such that $T_{V,\pi}M$ is the eigenbundle corresponding to the eigenvalue -1 of P while $T_{H,\pi}M$ is the eigenbundle corresponding to the eigenvalue $+1$ of P :

$$P \left(\frac{\partial}{\partial x^i} \right) = \frac{\partial}{\partial x^i} - 2N_i^\alpha \frac{\partial}{\partial y^\alpha}, \quad P \left(\frac{\partial}{\partial y^\alpha} \right) = -\frac{\partial}{\partial y^\alpha}. \tag{26}$$

So, our Corollary 4 can be thought in correspondence with Example 4 from [10, p. 894]; the present result is more efficient than Proposition 2.3 of [10, p. 891] since one of the generated almost Dirac structure is in fact already Dirac. The flatness of the connection is equivalent with the integrability of P expressed as the vanishing of the Nijenhuis tensor of P and then we recover the Theorem 3.3 of [10, p. 892].

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